

Effective field theory interpretation of CMS data

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Indirect searches for physics beyond the Standard Model

SM Effective Field Theory

One way to do this: The Standard Model Effective Field Theory (SMEFT)



- A consistent and model independent way to parametrise deviations in all SM processes

- There are several phenomena the standard model of particle physics (SM) can not explain, such as gravity, dark matter, neutrino masses, ...
- \rightarrow we know there must be *some* unknown theory beyond the standard model (BSM) that can explain these observations
- The LHC at CERN allows us to search for new particles with masses up to a few TeV, but none have been found so far
- Maybe BSM particles are just too heavy to be produced at the LHC? 🦻





 \rightarrow look for indirect evidence of BSM physics via deviations in precision measurements of known SM processes

• Constraints on Wilson coefficients c_i can then be matched to limits on parameters in BSM theories

1) Measurement

- Measure differential cross section of a process that is potentially sensitive to BSM effects
- Compare the measurement to theoretical predictions (assuming the SM)
- For example, pp $\rightarrow W\gamma$ (sensitive to modified triple gauge coupling: $\mathcal{Q}_W = \varepsilon^{ijk} W^{i\nu}_{\mu} W^{j\rho}_{\nu} W^{k\mu}_{\rho}$



2) Parameterization

• SMEFT operators modify the SM cross section

$$\sigma_{\text{total}} = \sigma_{\text{SM}} + \sum_{j} \frac{c_j}{\Lambda^2} \sigma_j^{\text{int.}} + \sum_{j,k} \frac{c_j c_k}{\Lambda^4} \sigma_{jk}^{\text{BSM}}$$

• The terms $\sigma_j^{\text{int.}}$ and σ_{jk}^{BSM} are computed using simulated data (MG5_aMC@NLO + SMEFTsim3)



3) Fit

- Construct a likelihood model based on the number of observed and predicted events in each bin (1), their uncertainties, and the SMEFT parameterization (2)
- $L = \prod \operatorname{Pois}(n_i | s_i(\vec{c}, \vec{\nu}) + b_i(\vec{\nu})) \prod p_k(\hat{\nu}_k | \nu_k)$
 - $-\vec{c}$: Wilson coefficients; $\vec{\nu}$: nuisance parameters corresponding to systematic uncertainties $-n_i$: observed number events in bin i $-s_i, b_i$: predicted number of signal and background events in bin *i*, for given values of \vec{c} and $\vec{\nu}$
- Run a maximum likelihood fit and determine the best fit values of the Wilson coefficients and their confidence intervals

 p_{τ}^{γ} (GeV)

CMS SMEFT Combination

- We combine seven sets of CMS measurements and electroweak precision observables (EWPO) measured at LEP and SLC
- Higgs sector: $H \rightarrow \gamma \gamma$
- Top sector: $t\overline{t}, t\overline{t}X$
- Electroweak sector: $W\gamma$, WW, $Z \rightarrow \nu\nu$, EWPO
- Strong sector: inclusive jet
- Inputs were selected to provide sensitivity to a broad set of SMEFT operators (64 in total), have negligible overlap in event selections, and small background contributions

Which input channel is sensitive to which operators:



Basis rotation

- Not enough data to constrain all 64 Wilson coefficients in a simultaneous fit (many are strongly correlated)
- Use Principal Component Analysis (PCA) to identify linear combinations of Wilson coefficients that can be constrained simultaneously



Future plans

- Add more existing measurements to combination (e.g. other Higgs decay modes, B-physics, ...)
- Dedicated analyses that target specific operators \rightarrow we recently started work on a triple-differential Drell-Yan analysis targetting 2-quark-2-lepton SMEFT operators (indirect sensitivity to Z' models)

Reference: CMS Collaboration, Combined effective field theory interpretation of Higgs boson, electroweak vector boson, top quark, and multi-jet measurements

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EV41 ($\lambda^{-1/2} = 4.4$) -0.1<mark>0.4 0.5</mark>-0.2<mark>-0.6</mark>-0.20.1 0.1 0.1-0.2 | | | <mark>| 0.3</mark>|0.1| |0.1| |0.1| <mark>|0.2</mark>| | | |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| |0.1| EV42 ($\lambda^{-1/2} = 4.9$)

 $\begin{bmatrix} c_1 & c_2 & c_3 & c_3 & c_3 & c_3 & c_3 & c_4 & c_5 & c$



